

# Stationary Phase Monte Carlo Methods

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## Introduction:

- Motivations and Outline of the Method;
  - phase transitions,
  - smoothing of the action,
- Theory Behind;
  - mollification/approximate identity method,
- Examples and Results;
  - Airy function action,
  - $\phi^4$  action,
- Conclusion.

## Outline:

The basic idea is to cast the partition function in such a way that its computation is manageable. Thus, the calculation of its derivatives (which yield the different Green's functions) become a feasible task, as well.

In order to do so, the full [complex] exponential weight is convoluted with a “probability-type” function, so that its oscillations are smoothed out and a Monte Carlo calculation can be performed.

This is particularly interesting when dealing with phase transitions since the factor  $\exp(iS(x))$  oscillates rapidly and these can be tamed — via the smoothing procedure that will be described shortly — so that a Monte Carlo computation can be performed. Also, coalescing stationary points are better handled by this method, again because of this smoothing technique.

## Theory:

The whole business is built upon the following construction:

$$\int f(x) dx \mapsto \int \langle f(x) \rangle_\epsilon dx , \quad (1)$$

where

$$\begin{aligned} \langle f(x) \rangle_\epsilon &\equiv (f * P_\epsilon)(x) , \\ &= \int P_\epsilon(x - y) f(y) dy , \end{aligned} \quad (2)$$

with

$$\int P_\epsilon(x) dx = 1 .$$

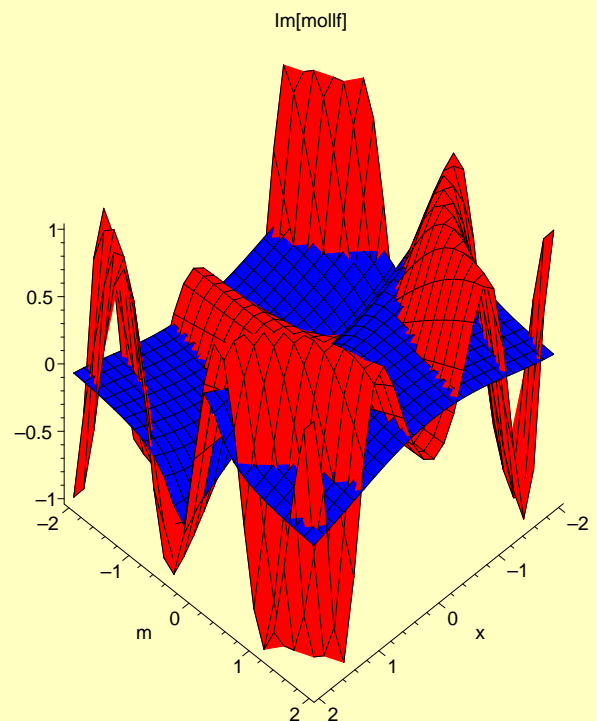
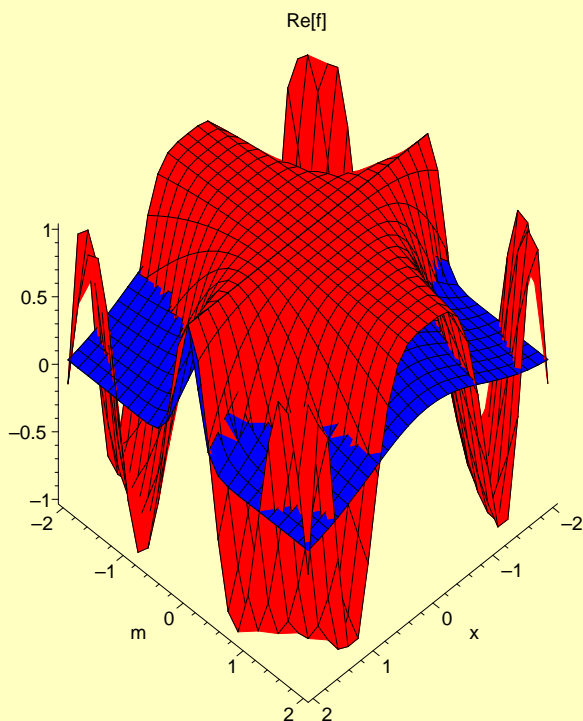
In this case, the smoothing function is given by  $P_\epsilon(x)$  which is called a “*mollifier*” or an “*approximate identity*”. The nice property that they have is:

$$\langle f(x) \rangle_\epsilon \xrightarrow{\epsilon \rightarrow 0} f(x) \Rightarrow \int \langle f(x) \rangle_\epsilon dx \xrightarrow{\epsilon \rightarrow 0} \int f(x) dx .$$

Simple example:

$$f_m(x) = \exp(i m x^2), \quad P_\epsilon(x) = \frac{\exp\left\{-\frac{1}{2} \frac{x^2}{\epsilon^2}\right\}}{\sqrt{2\pi\epsilon^2}}$$

$$\Rightarrow \langle f_m(x) \rangle_\epsilon = \frac{\exp\left\{\frac{i m x^2}{1-2 i m \epsilon^2}\right\}}{\sqrt{1-2 i m \epsilon^2}}$$



The method's application is in two steps:

1. Pre-averaging of the integrand: the generating functional is given by,

$$Z_\epsilon[j] = \frac{\int \langle e^{iS(\mathbf{x}) - i\mathbf{j}\cdot\mathbf{x}} \rangle_\epsilon [dx]}{\int \langle e^{iS(\mathbf{x})} \rangle_\epsilon [dx]} .$$

Thus, performing a *saddle-point expansion*:

$$P_\epsilon(\mathbf{y}) = \frac{\exp\left\{-\frac{1}{2}\mathbf{y}^\top \cdot (\epsilon^2)^{-1} \cdot \mathbf{y}\right\}}{\sqrt{(2\pi)^n \det(\epsilon^2)}} ,$$

$$\mathcal{J}_\epsilon = \int \langle f(\mathbf{x}) e^{iS(\mathbf{x}) - i\mathbf{j}\cdot\mathbf{x}} \rangle_\epsilon [dx]$$

$$\begin{aligned} &\approx f(\mathbf{x}_0) e^{iS(\mathbf{x}_0) - i\mathbf{j}\cdot\mathbf{x}_0} \times \\ &\times \int \frac{\exp\left\{-\frac{1}{2}\mathbf{B}^\top \cdot (\mathbb{1} + \epsilon^\top \cdot S'' \cdot \epsilon)^{-1} \cdot \mathbf{B}\right\}}{\sqrt{\det(\mathbb{1} + \epsilon^\top \cdot S'' \cdot \epsilon)}} [dx] , \end{aligned} \quad (3)$$

where  $\epsilon^2$  is the covariance matrix that defines the Gaussian distribution,  $\mathbf{B} = \epsilon \cdot \nabla S(\mathbf{x})$  and  $[S'']_{ij} = \partial^2 S(\mathbf{x}) / \partial x_i \partial x_j$  (Hessian matrix for  $S$ ).

2. Importance Sampling (MC): better handle the MC simulations. The equations become:

$$Z_\epsilon[j] = \frac{\int \left[ \frac{\langle e^{iS(\mathbf{x}) - i\mathbf{j}\cdot\mathbf{x}} \rangle_\epsilon}{W(\mathbf{x})} \right] [dW(\mathbf{x})]}{\int \left[ \frac{\langle e^{iS(\mathbf{x})} \rangle_\epsilon}{W(\mathbf{x})} \right] [dW(\mathbf{x})]} , \quad (4)$$

where  $W(\mathbf{x})$  is known as the *importance function* and the MC (Metropolis) sampling is done over this new measure. A reasonable choice for it is given by,

$$W(\mathbf{x}) = W_\epsilon(\mathbf{x}) = \left| \langle e^{iS(\mathbf{x})} \rangle_\epsilon \right| .$$

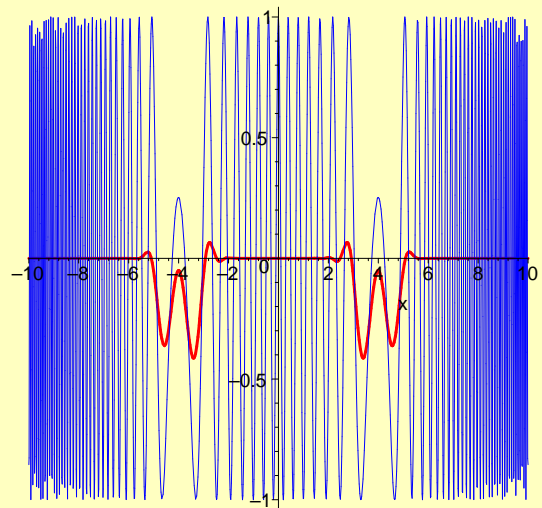
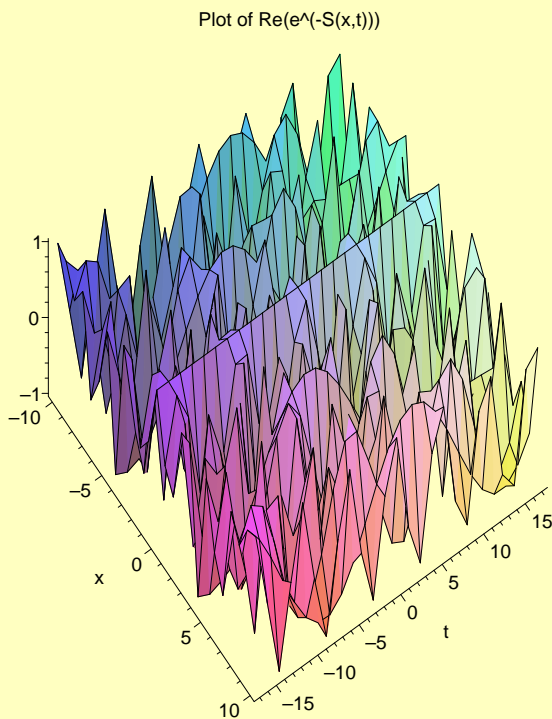
And, just as done before (saddle-point expansion), this becomes:

$$W_\epsilon(\mathbf{x}) = \left| \frac{\exp \left\{ iS(\mathbf{x}_0) - \frac{1}{2} \mathbf{B}^\top \cdot (\mathbf{1} + \epsilon^\top \cdot S'' \cdot \epsilon)^{-1} \cdot \mathbf{B} \right\}}{\sqrt{\det(\mathbf{1} + \epsilon^\top \cdot S'' \cdot \epsilon)}} \right|$$

## Example: Airy Function

The action is:  $S(x) = x^3/3 + tx$ . Thus,

$$\mathcal{Z}[t] = \frac{\int_{-\infty}^{\infty} \exp\left\{i \frac{x^3}{3} + i t x\right\} dx}{\int_{-\infty}^{\infty} \exp\left\{i \frac{x^3}{3}\right\} dx} \equiv \frac{\text{Ai}(t)}{\text{Ai}(0)}$$

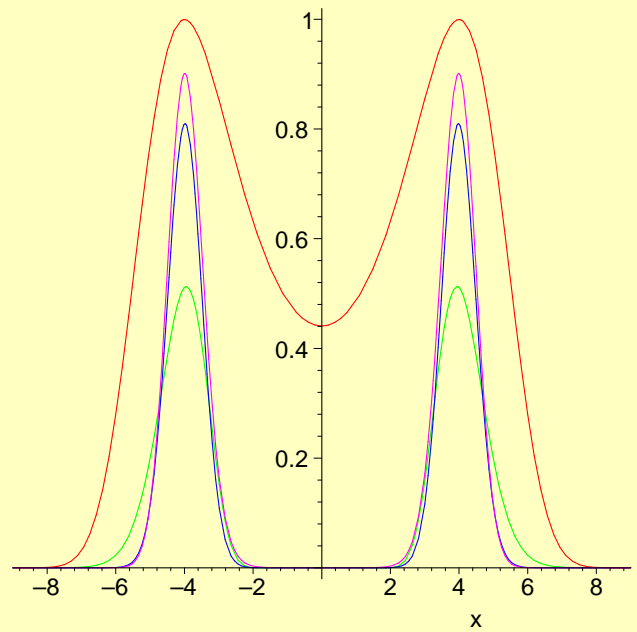
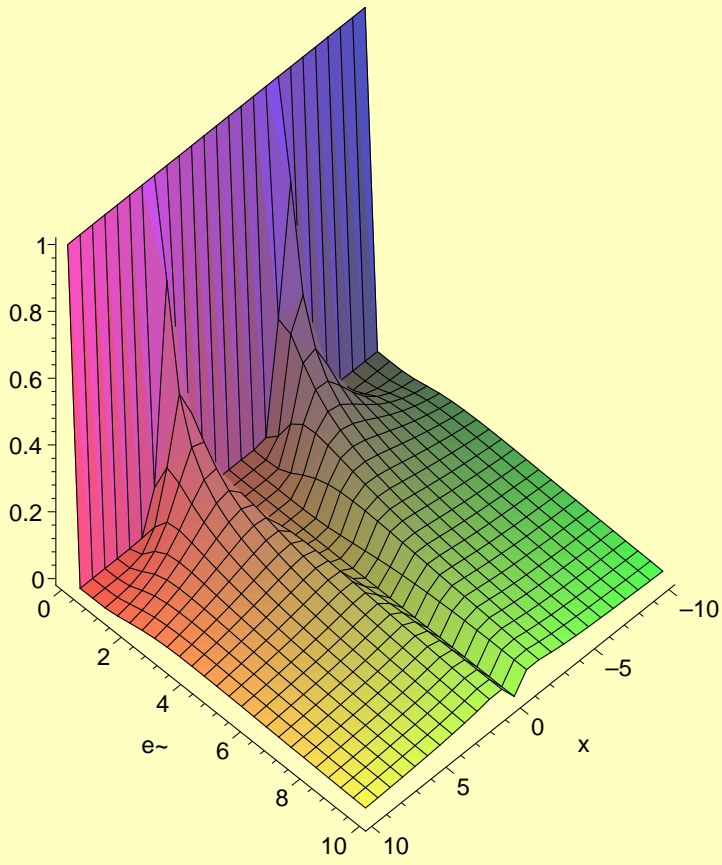


Legend

- $\text{Re}[\exp(-I^*s(x,-16))]$
- $\text{Re}[\exp(-I^*s(y,-16))]^4 P(x-y, 0.30)$



Plot of  $W(x, -16, e)$



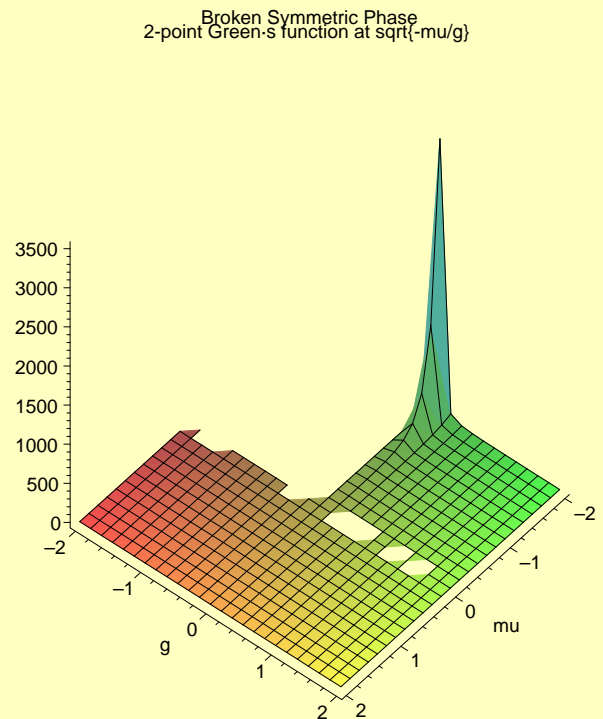
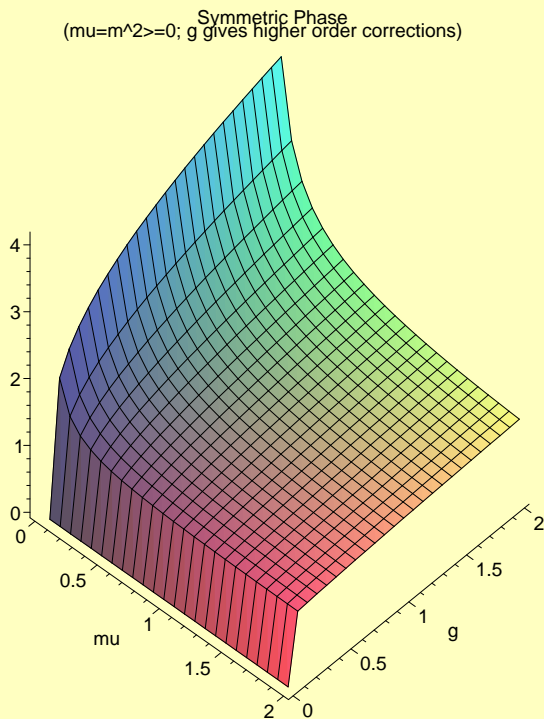
Legend

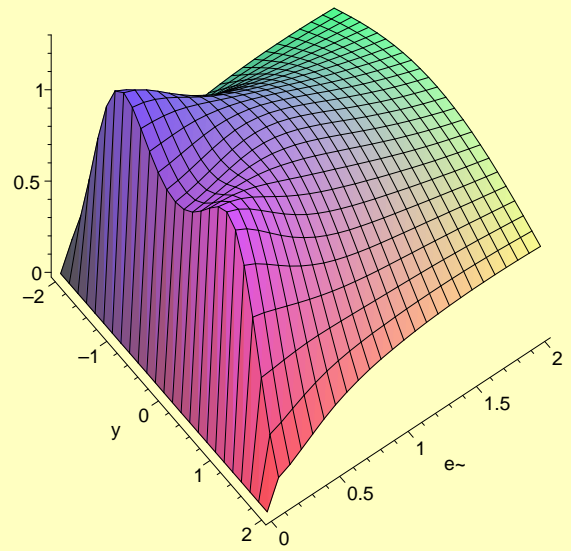
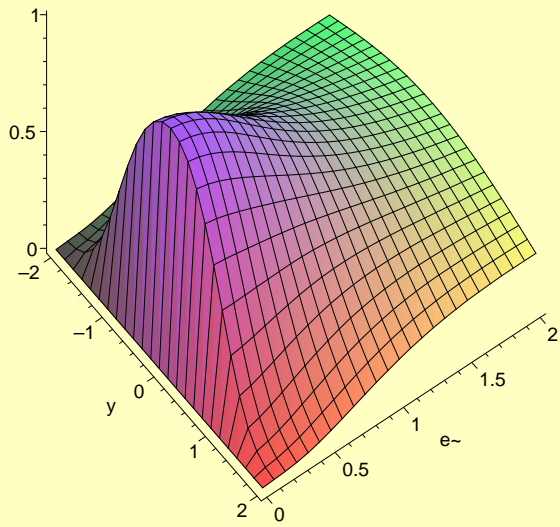
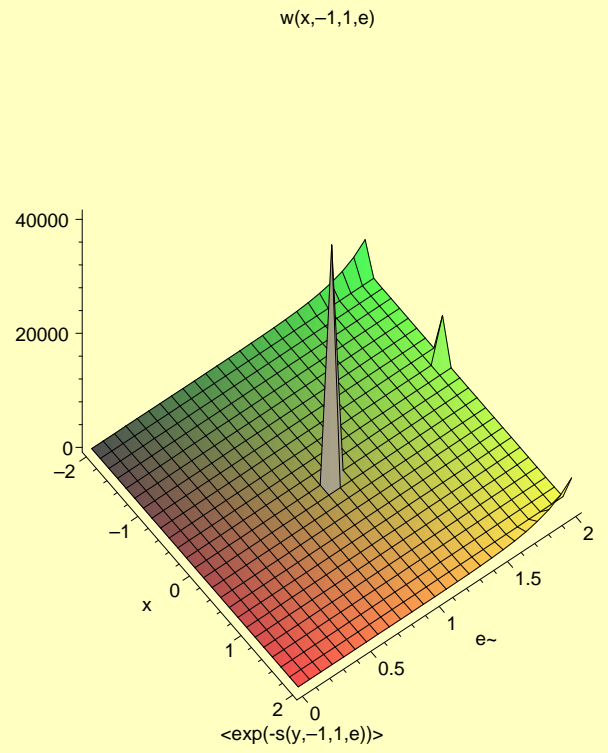
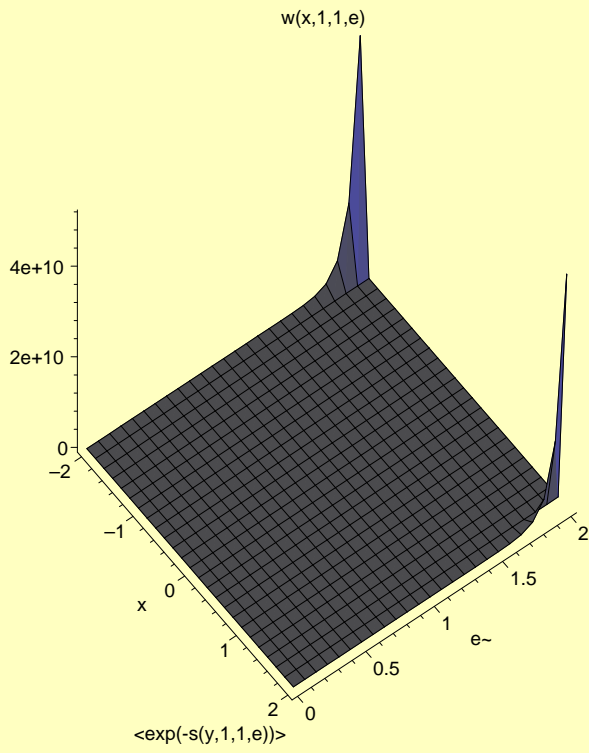
- $W(x, -16, 0.30)$
- $W(x, -16, 0.08)$
- $W(x, -16, 0.38)$
- $W(x, -16, 0.68)$

## Example: 0-dimensional $\phi^4$ Theory

The action is:  $S(\phi) = \mu \phi^2/2 + g \phi^4/4$ . Thus,

$$\mathcal{Z}_\epsilon[j] = \frac{\int_{-\infty}^{\infty} \left\langle \exp \left\{ i \frac{\mu}{2} \phi^2 + i \frac{g}{4} \phi^4 - i j \phi \right\} \right\rangle_\epsilon d\phi}{\int_{-\infty}^{\infty} \left\langle \exp \left\{ i \frac{\mu}{2} \phi^2 + i \frac{g}{4} \phi^4 \right\} \right\rangle_\epsilon d\phi}$$





## Conclusion:

This process can be interpreted as a way to implement block variables straight into the integrand, by means of smoothing it out. One big open question, under this view (coarse-graining, the graining parameter being  $\epsilon$ ), regards the renormalization procedure.

*We have just begun to calculate!*